

Stochastic transport equation and non-Lipshitz SDEs

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The linear transport equation (classically)

Given $b : \mathbb{R}_+ \times \mathbb{R}^d \rightarrow \mathbb{R}^d$ smooth vectorfield, \bar{u} smooth. Consider the Cauchy problem in $\mathbb{R}_+ \times \mathbb{R}^d$

$$\begin{cases} \partial_t u(t, x) + b(t, x) \cdot \nabla u(t, x) = 0 \\ u(0, x) = \bar{u}(x) \end{cases} \quad (1)$$

and the flow generated by b :

$$\begin{cases} \partial_t \Phi_{s,t}(x) = b(t, \Phi_{s,t}(x)) \\ \Phi_{s,s}(x) = x \end{cases}$$

Solutions to (1) are constant on the trajectories of b :

$$\frac{d}{dt} u(t, \Phi_{0,t}(x)) = \partial_t u(t, \Phi_{0,t}(x)) + \partial_t \Phi_{0,t}(x) \cdot \nabla u(t, \Phi_{0,t}(x)) = 0$$

Method of characteristics

The unique solution to (1) is $u(t, x) = \bar{u}(\Phi_{0,t}^{-1}(x))$.

Non-smooth vectorfields

Weak formulation

$$\begin{cases} \partial_t u + \operatorname{div}(bu) - (\operatorname{div} b)u = 0 \\ u(0, x) = \bar{u}(x) \end{cases}$$

Testing with smooth θ

$$\begin{aligned} \int \theta(x)u(t, x)dx &= \int \theta(x)\bar{u}(x)dx \\ &+ \int_0^t ds \int (u(s, x)b(s, x) \cdot \nabla\theta(x) + u(s, x)\theta(x)\operatorname{div} b(s, x))dx \end{aligned}$$

- ▶ Existence of L^∞ weak solutions when $b \in L^p$, $\operatorname{div} b \in L^1_{\text{loc}}$ and $\bar{u} \in L^\infty$
- ▶ **[DiPerna-Lions]** Renormalized solutions: uniqueness and stability of L^∞ weak solutions when $b \in L^1(W^{1,p}) \cap L^\infty$ and $\operatorname{div} b \in L^\infty$
- ▶ **[Ambrosio]** Renormalized solutions for **BV** vectorfields
- ▶ Use the transport equation to select a flow Φ defined *almost everywhere*

SDEs with non-smooth coefficients

Idea:

Perturb the equation of characteristics by an additive Brownian noise acting on all components.

Why?

Consider the SDE in \mathbb{R}^d

$$dX_t = b(t, X_t)dt + dW_t, \quad X_0 = x_0$$

- ▶ Strong solutions for b Lipschitz (+ linear growth) by fixed point method
- ▶ **[Veretennikov]** b bounded \Rightarrow uniqueness of strong solutions
- ▶ **[Krylov-Röckner]** Strong uniqueness for b in Sobolev spaces
- ▶ **[Davie]** b bounded \Rightarrow unique solution for a.e. Brownian path

\Rightarrow *The noise regularizes the flow of the vectorfield b* \Leftarrow

Stochastic flow

To implement the method of characteristics we need information on *dependence on initial conditions*.

Definition

A *stochastic flow* is a family of maps $\{\Phi_{s,t} : \mathbb{R}^d \rightarrow \mathbb{R}^d\}_{0 \leq s \leq t \leq T}$ such that

- ▶ $\Phi_{s,t}(x)$ is $\sigma(\{W_r - W_q\}_{s \leq q \leq r \leq t})$ measurable for any $x \in \mathbb{R}^d$, $0 \leq s \leq t \leq T$;
- ▶ $\lim_{t \rightarrow s^+} \Phi_{s,t}(x) = x$, a.s. for any x, s, t ;
- ▶ $\Phi_{u,t}(\Phi_{s,u}(x)) = \Phi_{s,t}(x)$

Theorem (Kunita)

If $b \in C^{1,\alpha}$ then there exists a $C^{1,\alpha'}$ -stochastic flow $\Phi_{s,t}$ for any $\alpha' < \alpha$ solving the SDE

$$\Phi_{s,t}(x) = x + \int_s^t b(u, \Phi_{s,u}(x)) du + W_t - W_s$$

for any $x \in \mathbb{R}^d$.

The Itô trick (I)

The regularization effect can be understood easily in the case $b(t, x) = b(x)$. Consider

$$X_t = x + \int_0^t b(X_s) ds + W_t$$

Try the *Itô trick*: interpret the integral over time as a correction in an Itô formula:

$$G(X_t) = G(x) + \int_0^t \nabla G(X_s) dW_s + \int_0^t LG(X_s) ds$$

with $L = \Delta/2 + b \cdot \nabla$. Assume that we can solve the elliptic problem

$$\lambda G - LG = b$$

for some $\lambda > 0$ (maybe very large), then

$$X_t + G(X_t) = x + G(x) + W_t + \int_0^t \nabla G(X_s) dW_s - \int_0^t \lambda G(X_s) ds$$

where G "has two derivatives more" than b . Setting $\psi(x) = x + G(x)$ we get

$$\psi(X_t) = \psi(x) + \int_0^t \nabla \psi(X_s) dW_s - \int_0^t \lambda G(X_s) ds$$

The Itô trick (II)

Theorem (Elliptic estimates)

For any $\epsilon > 0$, $\epsilon' < \epsilon$, $b \in C^\epsilon$, the elliptic equation $\lambda G - LG = b$ has a solution $G \in C^{2,\epsilon}$ for which $\|G\|_{2,\epsilon'} \rightarrow 0$ as $\lambda \rightarrow \infty$.

For λ large enough $\nabla\psi = 1 + \nabla G$ is invertible and ψ has inverse ψ^{-1} .
Let $Y_t = \psi(X_t)$, $y = \psi(x)$:

$$Y_t = y + \int_0^t \tilde{\sigma}(Y_s) dW_s + \int_0^t \tilde{b}(Y_s) ds$$

where $\tilde{\sigma}(y) = \nabla\psi \circ \psi^{-1}(y)$ and $\tilde{b}(y) = \lambda G \circ \psi^{-1}(y)$.

We have $\tilde{\sigma} \in C^{1,\epsilon'}$, $\tilde{b} \in C^{2,\epsilon'}$ and there exists a $C^{1,\epsilon'}$ -stochastic flow φ solving

$$\varphi_{s,t}(y) = y + \int_s^t \tilde{\sigma}(\varphi_{s,u}(y)) dW_u + \int_s^t \tilde{b}(\varphi_{s,u}(y)) du$$

Stochastic flow for C^ϵ vectorfields

By letting $\phi_{s,t} = \psi^{-1} \circ \varphi_{s,t} \circ \psi$ we obtain a $C^{1,\epsilon'}$ stochastic flow satisfying

$$\phi_{s,t}(x) = x + \int_s^t b(\phi_{s,u}(x)) du + W_t - W_s$$

- ▶ this flow is the unique strong solution to the SDE
- ▶ it does not depend on the choice of λ .
- ▶ we have an equation for $\nabla \phi_{s,t}(x)$:

$$\begin{aligned} \nabla \psi(\phi_{s,t}(x)) \nabla \phi_{s,t}(x) &= \nabla \psi(x) + \int_s^t \lambda \nabla G(\phi_{s,u}(x)) \nabla \phi_{s,u}(x) du \\ &\quad + \int_s^t \nabla^2 \psi(\phi_{s,u}(x)) \nabla \phi_{s,u}(x) dW_u \end{aligned}$$

- ▶ by a stopping procedure we can assume b locally in C^ϵ (+ linear growth)

Push-forward

For smooth b we have

$$\int \theta(\phi_{s,t}(x)) dx = \int \theta(x) \frac{dx}{J_{s,t}(x)}$$

where $J_{s,t}(x) = |\det \nabla \phi_{s,t}(x)|$ (Jacobian determinant) satisfy the differential equation

$$\frac{d}{dt} J_{s,t}(x) = \operatorname{div} b(\phi_{s,t}(x)) J_{s,t}(x), \quad J_{s,s}(x) = 1.$$

(the stochastic perturbation is solenoidal). Then

$$J_{s,t}(x) = \exp \left(\int_s^t \operatorname{div} b(\phi_{s,u}(x)) du \right)$$

For $b \in C^\epsilon$ by an approximation procedure and another Itô trick we get

$$J_{s,t}(x) = \exp \left(\Gamma(\phi_{s,t}(x)) - \Gamma(x) + \int_s^t \nabla \Gamma(\phi_{s,u}(x)) dW_u + \int_s^t \lambda \Gamma(\phi_{s,u}(x)) du \right)$$

where $\Gamma \in C^{1,\epsilon'}$ solve $\lambda \Gamma - L\Gamma = \operatorname{div} b$ in the sense of distributions.

Stochastic transport equation

The simplest stochastic perturbation which is compatible with the method of characteristics leads to the Stratonovich SPDE

$$\begin{cases} d_t u_t + b_t \cdot \nabla u_t dt + \sum_{i=1}^d \nabla_i u_t \circ dW_t^i = 0 \\ u_0(x) = \bar{u}(x) \end{cases}$$

and to the related SDE for the flow of characteristics:

$$\begin{cases} d_t \Phi_{s,t}(x) = b(t, \Phi_{s,t}(x)) dt + dW_t \\ \Phi_{s,s}(x) = x \end{cases}$$

Euristically we must have again $u_t(x) = \bar{u}(\Phi_{0,t}^{-1}(x))$.

Assume that b is locally bounded and $\operatorname{div} b \in L^q_{\text{loc}}$.

Definition

Given $\bar{u} \in L^p_{\text{loc}}$, for some $p \geq 1$ a solution of the stochastic transport equation (STE) in L^p_{loc} is a measurable function $(u(t, x, \omega), t \geq 0, x \in \mathbb{R}^d, \omega \in \Omega)$ such that

- (i) for P -a.e. $\omega \in \Omega, x \in \mathbb{R}^d, R > 0, \sup_{t \in [0, T]} \int_{B(x, R)} |u(t, x, \omega)|^p dx < \infty$
- (ii) for any test function $\theta \in C^0_0(\mathbb{R}^d)$, the process $t \mapsto \int_{\mathbb{R}^d} u(t, x)\theta(x)dx$ is continuous and \mathcal{F}_t -adapted;
- (iii) for any test function $\theta \in C^0_0(\mathbb{R}^d)$, the process $t \mapsto \int_{\mathbb{R}^d} u(t, x)\theta(x)dx$ is an \mathcal{F}_t -semimartingale satisfying

$$\begin{aligned} \int_{\mathbb{R}^d} u(t, x)\theta(x)dx &= \int_{\mathbb{R}^d} \bar{u}(x)\theta(x)dx + \sum_{i=1}^d \int_0^t \left(\int_{\mathbb{R}^d} u(s, x)D_i\theta(x)dx \right) \circ dW_s^i \\ &\quad + \int_0^t ds \int_{\mathbb{R}^d} u(s, x)[b(x) \cdot \nabla\theta(x) + \operatorname{div} b(x)\theta(x)]dx \end{aligned}$$

Main result

Theorem

Assume $b \in C^\epsilon$ and $\operatorname{div} b \in L^q$ and $\epsilon > d/q$. The STE has a unique solution u for any $\bar{u} \in L^p_{\text{loc}}$ and $u(t, x) = \bar{u}(\Phi_{0,t}^{-1}(x))$.

Note that by the pushforward formula

$$\int_{\mathbb{R}^d} f(x)g \circ \Phi_{s,t}(x)J_{s,t}(x)dx = \int_{\mathbb{R}^d} f \circ \Phi_{s,t}^{-1}(x)g(x)dx$$

with $J_{s,t}(x) \leq C$ locally. So if $f \in L^p_{\text{loc}}$, $g \in L^q_{\text{loc}}$ we have $f \circ \Phi_{s,t}^{-1} \in L^p_{\text{loc}}$ and

$$\int_A |f \circ \Phi_{s,t}^{-1}(x)|^p dx = \int_{\Phi_{s,t}^{-1}(A)} |f(x)|^p J_{s,t}(x) dx < \infty.$$

Existence

First we need to prove that $\int u(t, x)\theta(x)dx$ is a semimartingale.

Let $\phi_t = \phi_{0,t}$. Take a smooth test function θ , by Itô formula

$$\theta(\phi_t(y)) = \theta(y) + \int_0^t L^b \theta(\phi_s(y)) ds + \int_0^t \nabla \theta(\phi_s(y)) \cdot dW_s.$$

Let $J_t^\varepsilon(y)$ the Jacobian determinant of the flow ϕ_t^ε for the regularized vectorfield b^ε . Since b^ε is smooth: $dJ_t^\varepsilon(y) = \operatorname{div} b^\varepsilon(\phi_t(y))J_t^\varepsilon(y)dt$.

Then

$$\begin{aligned} \int_{\mathbb{R}^d} u_0(y)\theta(\phi_t(y))J_t^\varepsilon(y)dy &= \int_{\mathbb{R}^d} u_0(y)\theta(y)dy + \int_0^t ds \int_{\mathbb{R}^d} u_0(y)L^b \theta(\phi_s(y))J_s^\varepsilon(y)dy \\ &\quad + \int_0^t ds \int_{\mathbb{R}^d} u_0(y)\theta(\phi_s(y))\operatorname{div} b^\varepsilon(\phi_s(y))J_s^\varepsilon(y)dy \\ &\quad + \int_0^t dW_s \cdot \int_{\mathbb{R}^d} u_0(y)\nabla \theta(\phi_s(y))J_s^\varepsilon(y)dy \end{aligned}$$

In the limit $\varepsilon \rightarrow 0$ each term converges so

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} u_0(y)\theta(\phi_t(y))J_t^\varepsilon(y)dy = \int_{\mathbb{R}^d} u_0(y)\theta(\phi_t(y))J_t(y)dy = \int_{\mathbb{R}^d} u(t, y)\theta(y)dy$$

is a semi-martingale.

Next we need to prove that the semimartingale $\int u(t, x)\theta(x)dx$ satisfy the stochastic transport equation.

By the Stratonovic-Itô formula

$$\theta(\phi_t(y)) = \theta(y) + \int_0^t b \cdot \nabla\theta(\phi_s(y))ds + \int_0^t \nabla\theta(\phi_s(y)) \circ dW_s.$$

Then

$$\begin{aligned} \int_{\mathbb{R}^d} u_0(y)\theta(\phi_t(y))J_t^\varepsilon(y)dy &= \int_{\mathbb{R}^d} u_0(y)\theta(y)dy + \int_0^t ds \int_{\mathbb{R}^d} u_0(y)b \cdot \nabla\theta(\phi_s(y))J_s^\varepsilon(y)dy \\ &\quad + \int_0^t ds \int_{\mathbb{R}^d} u_0(y)\theta(\phi_s(y))\operatorname{div} b^\varepsilon(\phi_s(y))J_s^\varepsilon(y)dy \\ &\quad + \int_0^t dW_s \circ \int_{\mathbb{R}^d} u_0(y)\nabla\theta(\phi_s(y))J_s^\varepsilon(y)dy \end{aligned}$$

and take the limit $\varepsilon \rightarrow 0$ to conclude.

Uniqueness

Goal

Prove that, if $u(t, x)$ solve the STE then we must have $u(t, x) = \bar{u}(\phi_t^{-1}(x))$.

We start by smoothing u . Define

$$u_\varepsilon(t, y) = \int_{\mathbb{R}^d} u(t, x) \vartheta_\varepsilon(y - x) dx, \quad u_{0, \varepsilon}(y) = \int_{\mathbb{R}^d} u_0(x) \vartheta_\varepsilon(y - x) dx.$$

Since u is a solution to STE we get

$$\begin{aligned} u_\varepsilon(t, y) &= u_{0, \varepsilon}(y) + \int_0^t \left[\int_{\mathbb{R}^d} u(s, x) b(x) \cdot \nabla_x \vartheta_\varepsilon(y - x) dx \right] ds \\ &\quad + \int_0^t ds \int_{\mathbb{R}^d} u(s, x) \operatorname{div} b(x) \vartheta_\varepsilon(y - x) dx \\ &\quad + \sum_{i=1}^d \int_0^t \left[\int_{\mathbb{R}^d} u(s, x) D_{x_i} \vartheta_\varepsilon(y - x) dx \right] \circ dW_s^i \end{aligned}$$

Let $b^\delta = \vartheta_\delta * b$ and let ϕ^δ the associated flow.

By Stratonovich version of Itô-Wentzel calculus

$$\frac{d}{dt} u_\varepsilon(t, \phi_t^\delta(x)) = \left\{ \int u(t, x') \left[(b(x') - b^\delta(y)) \cdot \nabla_{x'} \vartheta_\varepsilon(y - x') + \operatorname{div} b(x') \vartheta_\varepsilon(y - x') \right] dx' \right\}$$

Test against $\rho \in C_0^\infty(\mathbb{R}^d)$ and perform a change of variables

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^d} u_\varepsilon(t, \phi_t^\delta x) \rho(x) dx \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} u(t, x') \left[(b(x') - b^\delta(y)) \cdot \nabla_{x'} \vartheta_\varepsilon(y - x') + \operatorname{div} b(x') \vartheta_\varepsilon(y - x') \right]_{y=\phi_t^\delta(x)} dx' \rho(x) dx \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} u(t, x') \left[(b(x') - b^\delta(y)) \cdot \nabla_{x'} \vartheta_\varepsilon(y - x') + \operatorname{div} b(x') \vartheta_\varepsilon(y - x') \right] dx' \rho((\phi_t^\delta)^{-1}(y)) \end{aligned}$$

By an integration by parts this is equal to

$$\begin{aligned} &= \int_{\mathbb{R}^d} \left[\int_{\mathbb{R}^d} \vartheta_\varepsilon(y - x') \left[b(x') - b^\delta(y) \right] \cdot \nabla_y \left[\rho((\phi_t^\delta)^{-1}(y)) J_t^\delta(y) \right] dy \right] u(t, x') dx' \\ & \quad + \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \left[\operatorname{div} b(x') - \operatorname{div} b^\delta(y) \right] \vartheta_\varepsilon(y - x') \rho((\phi_t^\delta)^{-1}(y)) J_t^\delta(y) dy u(t, x') dx' \end{aligned}$$

We want to show that both contributions go to zero as $\varepsilon \rightarrow 0$ and $\delta \rightarrow 0$

First term

$$\begin{aligned} A^\delta &= \lim_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^d} \vartheta_\varepsilon(\mathbf{y} - \mathbf{x}') [b(\mathbf{x}') - b^\delta(\mathbf{y})] \cdot \nabla_{\mathbf{y}} [\rho((\Phi_t^\delta)^{-1} \mathbf{y}) J_t^\delta(\mathbf{y})] d\mathbf{y} \\ &= [b(\mathbf{x}') - b^\delta(\mathbf{x}')] \cdot \nabla_{\mathbf{x}'} [\rho((\Phi_t^\delta)^{-1}(\mathbf{x}')) J_t^\delta(\mathbf{x}')] \end{aligned}$$

We can prove that

$$|\nabla [\rho((\Phi_t^\delta)^{-1}(\cdot)) J_t^\delta(\cdot)]| \lesssim \delta^\beta$$

locally as $\delta \rightarrow 0$ for any $\beta < -d/q$. Moreover

$$|b - b^\delta| \lesssim \delta^\epsilon$$

so $|A_\delta| \lesssim \delta^{\epsilon+\beta} \rightarrow 0$ as soon as $\epsilon + \beta > 0$.

Second term

$$\begin{aligned} & \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} [\operatorname{div} b(x') - \operatorname{div} b^\delta(y)] \vartheta_\varepsilon(y - x') \rho((\Phi_t^\delta)^{-1}(y)) J_t^\delta(y) dy u(t, x') dx' \\ &= \int_{\mathbb{R}^d} \operatorname{div} b(x') \left(\int_{\mathbb{R}^d} \vartheta_\varepsilon(y - x') \rho((\Phi_t^\delta)^{-1}(y)) J_t^\delta(y) dy \right) u(t, x') dx' \\ &\quad - \int_{\mathbb{R}^d} \operatorname{div} b^\delta(y) \rho((\Phi_t^\delta)^{-1}(y)) J_t^\delta(y) u_\varepsilon(t, y) dy \end{aligned}$$

and both terms converge, as $\varepsilon \rightarrow 0$ followed by $\delta \rightarrow 0$ to

$$\int_{\mathbb{R}^d} \operatorname{div} b(y) \rho(\Phi_t^{-1}(y)) J_t(y) u(t, y) dy$$

so their difference converge to zero.

We obtained

$$\lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} \left[\int_{\mathbb{R}^d} u_\varepsilon(t, \phi_t^\delta x) \rho(x) dx - \int_{\mathbb{R}^d} u_\varepsilon(0, x) \rho(x) dx \right] = 0.$$

Now

$$\begin{aligned} \int_{\mathbb{R}^d} u_\varepsilon(t, \phi_t^\delta x) \rho(x) dx &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} u_\varepsilon(t, y) \vartheta_\varepsilon(\phi_t^\delta(x) - y) \rho(x) dx dy \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} u_\varepsilon(t, y) \vartheta_\varepsilon(z - y) \rho((\phi_t^\delta)^{-1}(z)) J_t^\delta((\phi_t^\delta)^{-1}(z))^{-1} dz dy \\ &\rightarrow \int_{\mathbb{R}^d} u(t, z) \rho(\phi_t^{-1}(z)) J_t(\phi_t^{-1}(z))^{-1} dz \end{aligned}$$

This yields

$$\int_{\mathbb{R}^d} u(t, z) \rho(\phi_t^{-1}(z)) J_t(\phi_t^{-1}(z))^{-1} dz = \int_{\mathbb{R}^d} u(0, x) \rho(x) dx$$

for every $\rho(x) \in C_0^\infty(\mathbb{R}^d)$. Choosing ρ appropriately we get

$$\int_{\mathbb{R}^d} u(t, z) \rho(z) dz = \int_{\mathbb{R}^d} u(0, x) \rho(\phi_t(x)) J_t(x) dx = \int_{\mathbb{R}^d} u(0, \phi_t^{-1}(y)) \rho(y) dy.$$

Counterexamples to certain extensions

Example (Random vectorfields)

Take $b(t, x) = \sqrt{|x - W_t|}$, then

$$dX_t = b(t, X_t)dt + dW_t = \sqrt{|X_t - W_t|}dt + dW_t.$$

By the change of variables $Y_t = X_t - W_t$ we obtain

$$dY_t = \sqrt{|Y_t|}dt$$

so path-wise uniqueness is impossible in general.

Not so artificial...

Consider a 2d stochastic Euler equation in vorticity variables

$$\partial_t \xi(t, x) + (u(t, x) \cdot \nabla \xi(t, x)) dt + \nabla \xi(t, x) \circ dW(t) = 0$$

where $\xi = \partial_2 u_1 - \partial_1 u_2$.

Formally equivalent to the "system" of stochastic ordinary equations

$$dX_t^a = \left[\int_{\mathbb{R}^2} K(X_t^a - X_t^{a'}) \xi_0(X_t^{a'}) da' \right] dt + dW_t, \quad a \in \mathbb{R}^2$$

for a suitable kernel K , ξ_0 being the initial condition of the vorticity equation.

By the change of variable $Y_t^a = X_t^a - W_t$ we obtain

$$dY_t^a = \left[\int_{\mathbb{R}^2} K(Y_t^a - Y_t^{a'}) \xi_0(X_t^{a'}) da' \right] dt$$

The equation for (Y_t^a) corresponds to the classical vorticity equation

$$\frac{\partial_t \xi'(t, x)}{\partial t} + (u'(t, x) \cdot \nabla \xi'(t, x)) dt = 0 \quad \xi' = \partial_2 u'_1 - \partial_1 u'_2$$

with initial condition ξ_0 .

Possible way out

Consider a more complex (infinite-dimensional) noise:

$$dX_t^a = \left[\int_{\mathbb{R}^2} K(X_t^a - X_t^{a'}) \xi_0(X_t^{a'}) da' \right] dt + \sum_{k=1}^{\infty} \sigma_k(X_t^a) dW_t^k, \quad a \in \mathbb{R}^2$$

where each point X_a is moved "almost" independently of the others.
Seems useful to require

$$\sum_{k=1}^{\infty} \sigma_k(x) \sigma_k(y) = a(|x - y|)$$

with $a(r) \simeq r^\alpha$ as $r \rightarrow 0$, $\alpha > 0$. This in order to hope some regularizing effect of the noise over the deterministic (and singular) drift.

Connection with the theory of stochastic flows of Le Jan-Raimond.