

Equations for Euclidean fields



prologue

the origins

Itô's original paper

(Japanese version 1942, M.A.M.S. 1951)

Differential Equations Determining a Markoff Process*

KIYOSI ITÔ

More generally, for a simple Markoff process with its states being represented by the real numbers and having continuous parameter, the problem of determining quantities corresponding to $p_{ij}^{(k)}$ mentioned above and of constructing the corresponding Markoff process once these quantities are given has been investigated systematically by Kolmogoroff[3], who reduced the problem to the study of differential equations or integro-differential equations satisfied by the transition probability function.

W. Feller[4] has proved under fairly strong assumptions that these equations possess a unique solution and furthermore that the solution exhibits the properties of transition probability function.

However, if we adopt more strict point of view such as the one J. Doob[5] has applied toward his investigation of stochastic processes, it seems to us that the aforementioned work done by Feller is not quite adequate. For example, even though the differential equation determining the transition probab-

ility function of a continuous stochastic process was solved in §3 of that paper, no proof was given of the fact that it is possible to introduce by means of this solution a probability measure on some "continuous" function space.

The objective of this article, then, is:

- 1) to formulate the problem precisely, and
- 2) to give a rigorous proof, à la Doob, for the existence of continuous parameter stochastic processes.

§1. Definition of Differentiation of a Markoff Process

Let $\{y_t\}$ be a (simple) Markoff process and denote by $F_{t_0,t}$ the conditional probability distribution¹ of $y_t - y_{t_0}$ given that " y_{t_0} is determined". $F_{t_0,t}$ is clearly a P_{y_t} -measurable (ρ) function² of y_t , where ρ denotes the Lévy distance among probability distributions.

Definition 1.1.³

If

$$(1) \quad F_{t_0,t}^{*[1/t-t_0]}$$

(here $[a]$ is the integer part of the number a , and " $*k$ " denotes the k -fold convolution) converges in probability with respect to the Lévy distance ρ as $t \rightarrow t_0 + 0$, then we call the limit random variable (taking values in the space of probability distributions) the derivative of $\{y_t\}$ at t_0 and denote it by

$$(2) \quad D_{t_0}\{y_t\} \text{ or } Dy_{t_0}.$$

Corollary 1.1. Dy_{t_0} is an infinitely divisible probability distribution.⁴

ability distribution.

Dy_{t_0} obtained above is a function of t_0 as well as of y_{t_0} , and so, we denote it by $L(t_0, y_{t_0})$ corresponds precisely to the "basic transition probability" discussed in the Introduction.

The precise formulation of the problem of Kolmogoroff, then, is to solve the equation

$$(4) \quad Dy_t = L(t, y_t)$$

when the quantity $L(t, y)$ is given.

§2. A Comparison Theorem

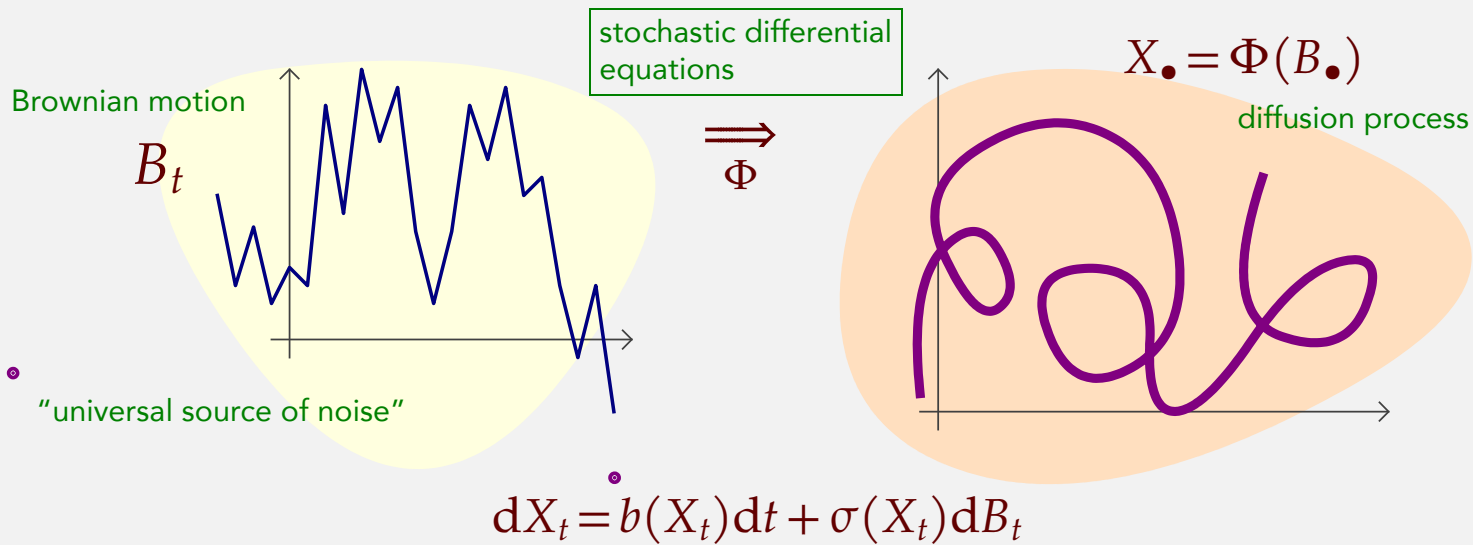
Let us prove a comparison theorem for Dy_{t_0} which we shall make use of later.

Theorem 2.1. *Let $\{y_t\}$, $\{z_t\}$ be simple Markoff processes satisfying the following conditions:*

- (1) $y_{t_0} = z_{t_0}$.
- (2) $E(y_t - z_t \mid y_{t_0}) = o(t - t_0)$, where o is the Landau symbol.
- (3) $\sigma(y_t - z_t \mid y_{t_0}) = o(\sqrt{t - t_0})$.

(Here $E(x \mid y)$ denotes the conditional expectation of x given y and $\sigma(x \mid y)$ denotes the conditional standard deviation of x given y . Also, the quantity o may depend on t_0 or y_{t_0}). Then, whenever Dz_{t_0} exists, Dy_{t_0} exists also, and $Dy_{t_0} = Dz_{t_0}$ holds.

Ito's brilliant idea



act I

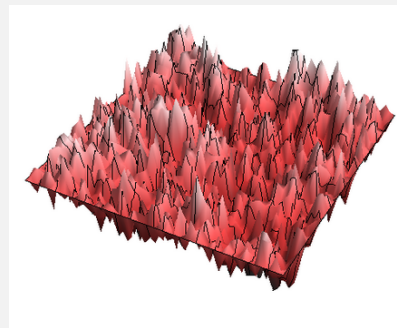
Markov fields in $d \geq 2$

Euclidean quantum fields

Feynman–Kac formula / non-perturbative path integral formulation of QFTs

$$\nu(d\varphi) = \frac{\exp[-S(\varphi)] \mathcal{D}\varphi}{Z}, \quad \varphi \in \mathcal{S}'(\mathbb{R}^d)$$

$$S(\varphi) = \frac{1}{2} \int_{\mathbb{R}^d} (|\nabla \varphi(x)|^2 + m^2 \varphi(x)^2) dx + \int_{\mathbb{R}^d} V(\varphi(x)) dx$$



$$V(\varphi) = \lambda \varphi^4 + \dots, \quad \lambda \cos(\beta \varphi), \quad \lambda \cosh(\beta \varphi), \quad \lambda \exp(\beta \varphi), \quad \dots, \quad d = 1, 2, 3$$

σ -models, gauge theories,...

(Markovian) probability measure ν on $\mathcal{S}'(\mathbb{R}^d)$ · natural generalisations of diffusion processes

Clay Millennium problem: prove existence of non-Abelian Euclidean quantum Yang–Mills theory in $d = 4$

Gaussian free field

▷ GFF · simplest example of EQFT · Gaussian measure μ on $\mathcal{S}'(\mathbb{R}^d)$ s.t.

$$\int \varphi(x)\varphi(y)\mu(d\varphi) = G(x-y) = \int_{\mathbb{R}^d} \frac{e^{ik(x-y)}}{m^2 + |k|^2} \frac{dk}{(2\pi)^d} = (m^2 - \Delta)^{-1}(x-y), \quad x, y \in \mathbb{R}^d$$

and zero mean · $m > 0$ is the *mass* · $G(0) = +\infty$ if $d \geq 2$: not a function · distribution of regularity

$$\alpha < (2-d)/2$$

▷ can be used to construct a QFT but the theory is free: no interaction

the standard recipe for non-Gaussian Euclidean fields

- ❶ go on a periodic lattice: $\mathbb{R}^d \rightarrow \mathbb{Z}_{\varepsilon, L}^d = (\varepsilon \mathbb{Z} / 2\pi L \mathbb{N})^d$ with spacing $\varepsilon > 0$ and side $2\pi L$

$$\int F(\varphi) v^{\varepsilon, L}(\mathrm{d}\varphi) = \frac{1}{Z_{\varepsilon, L}} \int_{\mathbb{R}^{\mathbb{Z}_{\varepsilon, L}^d}} F(\varphi) e^{-\frac{1}{2} \varepsilon^d \sum_{x \in \mathbb{Z}_{\varepsilon, L}^d} \overbrace{|\nabla_\varepsilon \varphi(x)|^2 + m^2 \varphi(x)^2 + V_\varepsilon(\varphi(x))}^{S_\varepsilon(\varphi)}} \mathrm{d}\varphi$$

ε is an UV regularisation and L the IR regularisation

- ❷ choose V_ε appropriately so that $v^{\varepsilon, L} \rightarrow v$ to some limit as $\varepsilon \rightarrow 0$ and $L \rightarrow \infty$. E.g. take V_ε polynomial bounded below. $d=2, 3$.

$$V_\varepsilon(\xi) = \lambda(\xi^4 - a_\varepsilon \xi^2)$$

The limit measure will depend on $\lambda > 0$ and on $(a_\varepsilon)_\varepsilon$ which has to be s.t. $a_\varepsilon \rightarrow +\infty$ as $\varepsilon \rightarrow 0$. It is called the Φ_d^4 measure.

- ❸ study the possible limit points [the Φ_d^4 measure] · ask interesting questions: uniqueness? non-uniqueness? decay of correlations? intrinsic description?

some models

▷ $d=1$ · time-reversal symmetric, translation invariant, Markov diffusions. the generator is given by an implicit expression involving the ground state Ψ of the Hamiltonian H

$$\mathcal{L} = \nabla \log \Psi \cdot \nabla + \frac{1}{2} \Delta \quad H = -\Delta + x^2 + V(x).$$

▷ $d=2$ · various choices ($a_\varepsilon \rightarrow +\infty$)

$$V_\varepsilon(\xi) = \lambda \xi^{2l} + \sum_{k=0}^{2l-1} a_{k,\varepsilon} \xi^k, \quad V_\varepsilon(\xi) = a_\varepsilon \cos(\beta \xi)$$

$$V_\varepsilon(\xi) = a_\varepsilon \cosh(\beta \xi), \quad V_\varepsilon(\xi) = a_\varepsilon \exp(\beta \xi)$$

▷ $d=3$ · “only” 4th order (6th order is critical)

▷ $d=4$ · all the possible limits are Gaussian (see Aizenmann–Duminil Copin)

act II

the quest for equations

11. Remarks on Markov field equations

E. NELSON

11.1. Introduction

WE have no new existence theorems in constructive quantum field theory to present here, but we wish to indicate a new direction which looks promising and which certainly poses many interesting questions.

Only the theory of a neutral scalar field with a quartic self-interaction will be considered. This theory has been much studied in dimension $d = 2$ and Glimm and Jaffe have pioneered the study in dimension $d = 3$. (For references, see [4]—in particular, see the first article by Glimm, Jaffe, and Spencer and the reference listed there.)

We wish to stress field equations, and so we will begin with a heuristic discussion from that point of view. The simplest non-linear relativistic field equation with good formal properties is

$$(\square + m^2)A = -gA^3 + \alpha A, \quad (11.1)$$

corresponding to the interaction Lagrangian density $-(g/4)A^4 + (\alpha/2)A^2$. We could of course absorb the linear term αA in the term $m^2 A$, but we prefer not to. Here m^2 and g are positive.

The Euclidean approach to the problem of quantized solutions to (11.1) is, in rough outline, as follows. The Wightman distributions (vacuum expectation values of products of field operators) are the boundary values of analytic functions which in the Euclidean region of points with real space

11.3. The Markov field equation

Let us compute:

$$\begin{aligned}
 \hat{E}\phi(x) - \mu &= \frac{1}{N} \int_{-\infty}^{\infty} (\xi - \mu) \exp\left[\left(-\frac{g}{4}\xi^4 + \frac{\alpha}{2}\xi^2\right)\varepsilon^d\right] \exp\left[-\frac{(\xi - \mu)^2}{2\sigma^2}\right] d\xi \\
 &= \frac{1}{N} \int_{-\infty}^{\infty} \exp\left[\left(-\frac{g}{4}\xi^2 + \frac{\alpha}{2}\xi^2\right)\varepsilon^d\right] \left(-\sigma^2 \frac{d}{d\xi} \exp\left[-\frac{(\xi - \mu)^2}{2\sigma^2}\right]\right) d\xi \\
 &= \frac{1}{N} \int_{-\infty}^{\infty} \left(\sigma^2 \frac{d}{d\xi} \exp\left[\left(-\frac{g}{4}\xi^4 + \frac{\alpha}{2}\xi^2\right)\varepsilon^d\right]\right) \exp\left[-\frac{(\xi - \mu)^2}{2\sigma^2}\right] d\xi \\
 &= \frac{1}{N} \int_{-\infty}^{\infty} \sigma^2 (-g\xi^3 + \alpha\xi) \varepsilon^d \exp\left[\left(-\frac{g}{4}\xi^4 + \frac{\alpha}{2}\xi^2\right)\varepsilon^d\right] \times \\
 &\quad \times \exp\left[-\frac{(\xi - \mu)^2}{2\sigma^2}\right] d\xi \\
 &= \sigma^2 \hat{E}[-g\phi(x)^3 + \alpha\phi(x)]\varepsilon^d.
 \end{aligned}$$

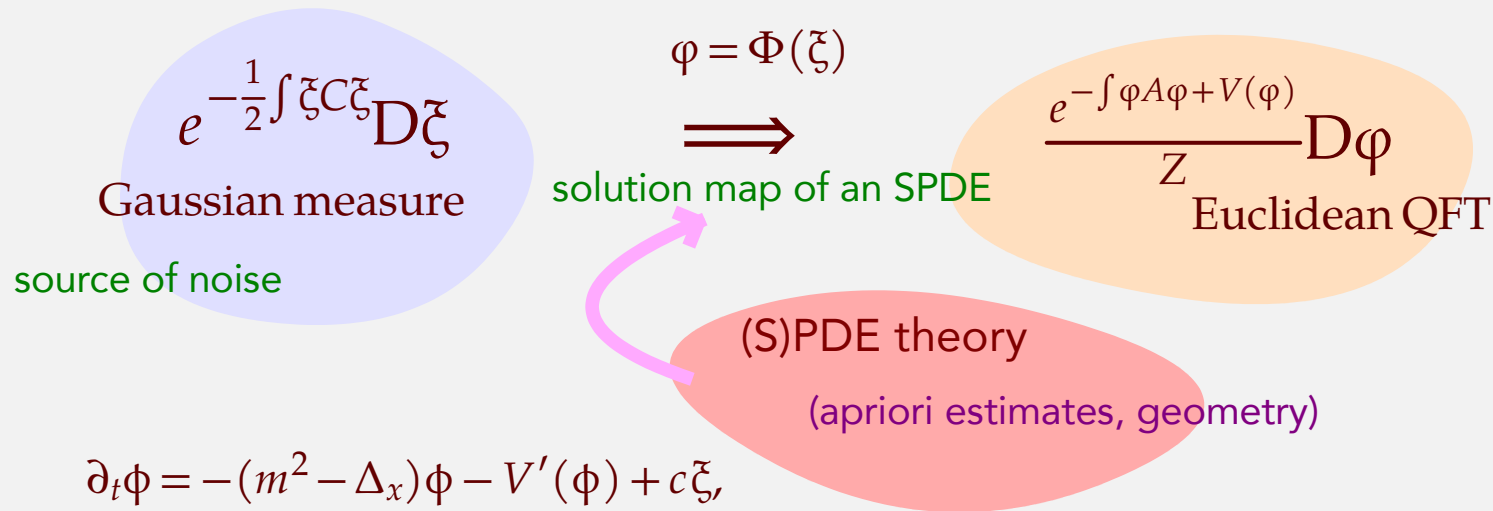
We may write this as

$$\phi(x) - \mu = \sigma^2 [-g\phi(x)^3 + \alpha\phi(x)]\varepsilon^d + \sigma^2 \omega(x)\varepsilon^d, \quad (11.8)$$

where $\omega(x)$ is a function of $\phi(x)$ and $\phi(y)$ for the nearest neighbours y of x (because (11.8) is the definition of $\omega(x)$) and

$$\hat{E}\omega(x) = 0. \quad (11.9)$$

stochastic quantisation map



systematic use of tools from **PDE theory**, recent advances in **singular PDEs**, **renormalization group** ideas coupled with pathwise analysis

stochastic equations for the free Gaussian free field

Gaussian free field $\mu : \mathbb{E}[\varphi(x)\varphi(y)] = (m^2 - \Delta)^{-1}(x - y) \cdot \xi$ white noise

❶ “Gaussian map”:

$$\varphi(x) = (m^2 - \Delta)^{-1/2} \xi(x), \quad (m^2 - \Delta)\varphi(x) = (m^2 - \Delta)^{1/2} \xi(x), \quad x \in \mathbb{R}^d$$

❷ Stochastic mechanics (Nelson):

$$\partial_{x_0} \varphi(x_0, \bar{x}) = -(m^2 - \Delta_{\bar{x}})^{1/2} \varphi(x_0, \bar{x}) + \xi(x_0, \bar{x}), \quad x_0 \in \mathbb{R}, \bar{x} \in \mathbb{R}^{d-1}$$

❸ Parabolic stochastic quantization (Parisi–Wu):

$$\varphi(x) \sim \phi(t, x) \quad \partial_t \phi(t, x) = -(m^2 - \Delta_x) \phi(t, x) + c \xi(t, x), \quad t \in \mathbb{R}, x \in \mathbb{R}^d$$

❹ Elliptic stochastic quantization (Parisi–Sourlas):

$$\varphi(x) \sim \phi(z, x) \quad (-\Delta_z) \phi(z, x) = -(m^2 - \Delta_x) \phi(z, x) + c \xi(z, x), \quad z \in \mathbb{R}^2, x \in \mathbb{R}^d$$

stochastic equations for non-Gaussian EQFTs ($V \neq 0$)

- ❶ Shifted Gaussian map (Albeverio/Yoshida) **[does not have the right properties!]**

$$(m^2 - \Delta)\varphi(x) + V'(\varphi(x)) = (m^2 - \Delta)^{1/2}\xi(x), \quad x \in \mathbb{R}^d$$

- ❷ Stochastic mechanics (Nelson): ground-state transformation **[implicit!]**

$$\partial_{x_0}\varphi(x_0, \bar{x}) = [\nabla_{\varphi(x_0, \bar{x})} \log \Psi(\varphi)] + \xi(x_0, \bar{x}), \quad x_0 \in \mathbb{R}, \bar{x} \in \mathbb{R}^{d-1}$$

- ❸ Parabolic stochastic quantization (Parisi–Wu): Langevin diffusion

$$\varphi(x) \sim \phi(t, x) \quad \partial_t \phi(t, x) = -(m^2 - \Delta_x)\phi(t, x) - V'(\phi(t, x)) + c\xi(t, x), \quad t \in \mathbb{R}, x \in \mathbb{R}^d$$

- ❹ Elliptic stochastic quantization (Parisi–Sourlas):

$$\varphi(x) \sim \phi(z, x) \quad (-\Delta_z)\phi(z, x) = -(m^2 - \Delta_x)\phi(z, x) - V'(\phi(z, x)) + c\xi(z, x), \quad z \in \mathbb{R}^2, x \in \mathbb{R}^d$$

an short history of (rigorous) stochastic quantisation for bosonic fields

- ▶ 1981 · Parisi/Wu – stochastic quantisation for gauge theories (SQ)
- ▶ 1985 · Jona-Lasinio/Mitter · “On the stochastic quantization of field theory” (rigorous SQ for Φ_2^4 on bounded domain)
- ▶ 1988 · Damgaard/Hüffel · review book on SQ (theoretical physics)
- ▶ 1990 · Funaki · Control of correlations via SQ (smooth reversible dynamics)
- ▶ 2003 · Da Prato/Debussche · “Strong solutions to the stochastic quantization equations”
- ▶ 2014 · Hairer – Regularity structures, local dynamics of Φ_3^4
- ▶ 2017 · Mourrat/Weber · global solutions for Φ_2^4 , coming down from infinity for Φ_3^4
- ▶ 2018 · Albeverio/Kusuoka · “The invariant measure and the flow associated to $\Phi_3^4 \dots$ ”
- ▶ 2021 · Hofmanova/G. · Global space-time solutions for Φ_3^4 + ~OS axioms + IbP formula
- ▶ 2022 · Hairer/Steele · “optimal” tail estimates
- ▶ 2020–2023 · Chandra/Chevyrev/Hairer/Shen · SQ for Yang–Mills 2d/3d (local theory)
- ▶ 2023 · Duch/G./Rinaldi · SQ for fractional Φ^4 models in the full subcritical regime + ~OS axioms
- ▶ 2024 · Bringmann/Cao · SQ for 2d Abelian-Higgs model (global-in-time theory on a torus)
- ▶ 2025 · Barashkov/Gunaratnam · Markov property for Φ_3^4 on cylinders
- ▶ 2025 · G./Meyer · Sine-Gordon up to 6π via FBSDEs + ~OS axioms
- ▶ 2025 · Duch/Hairer/Yi/Zhao · Uniqueness in high temperature for $\Phi_{2,3}^4$
- ▶ 2025 · Bauerschmidt/Dagallier/Weber · Uniqueness up to critical temperature for Φ_2^4 (Holley–Stroock)

act III

challenges for the stochastic analysts

random variables are not enough

Let φ be the GFF in $d \geq 3$, then the Wick powers $[\![\varphi^n]\!]$ are **not** well defined random distributions!

$$\mathbb{E} \left[[\![\varphi^n(x)]\!] [\![\varphi^n(y)]\!] \right] = n! \left(\mathbb{E} [\varphi(x) \varphi(y)] \right)^n \approx |x - y|^{(2-d)n} \notin L^1_{\text{loc}}(\mathbb{R}_x^d \times \mathbb{R}_y^d)$$

So the natural local fields $[\![\varphi^n]\!]$ (a.k.a. **composite fields**) does not have a place in the classical probabilistic setting.

☞ For Gaussian fields one can use the Fock space picture (or chaos decomposition) to define generalized random variables, i.e. elements in the dual of sufficiently smooth (in the sense of Malliavin) random variables.

It is not clear how to generalise this to non-Gaussian fields, especially those which are not absolutely continuous wrt. the GFF.

Composite fields appears naturally in integration by parts formulas.

integration by parts formulas

Q: how to characterize an Euclidean field?

a physically motivated approach · IbP formulas (see Schwinger–Dyson equations)

$$\int_{\mathcal{S}'(\mathbb{R}^d)} \left[\frac{\delta}{\delta\varphi} - \frac{\delta S}{\delta\varphi}(\varphi) \right] F(\varphi) \nu(d\varphi) = 0, \quad \forall F \text{ in some nice class}$$

$$\left[\frac{\delta}{\delta\varphi} - \frac{\delta S}{\delta\varphi}(\varphi) \right] \approx e^{S(\varphi)} \frac{\delta}{\delta\varphi} e^{-S(\varphi)}$$

▷ existence of solutions, uniqueness? Note: related to generator of the Langevin dynamics

$$\left[\frac{\delta}{\delta\varphi} - \frac{\delta S}{\delta\varphi}(\varphi) \right] \frac{\delta F(\varphi)}{\delta\varphi} = \mathcal{L}F.$$

very little is known mathematically. [(in finite dim) V. I. Bogachev, N. V. Krylov, and M. Röckner · Elliptic and Parabolic Equations for Measures · *Russian Mathematical Surveys* (2009)]

Euclidean fields add substantial problems. [A. I. Kirillov · On the Reconstruction of Measures from Their Logarithmic Derivatives · *Izvestiya: Mathematics* (1995)]

differential characterization of the $\exp(\beta\varphi)_2$ model

[F. C. De Vecchi, MG and M. Turra · A Singular Integration by Parts Formula for the Exponential Euclidean QFT on the Plane · arXiv (2022)]

Heuristically

$$\int_{\mathcal{S}'(\mathbb{R}^d)} \left[\frac{\delta}{\delta\varphi} - (m^2 - \Delta + \lambda \llbracket \exp(\beta\varphi) \rrbracket) \right] F(\varphi) \nu(d\varphi) = 0, \quad \forall F \in \mathcal{C}$$

Rigorously

$$\lim_{\varepsilon \rightarrow 0} \int_{\mathcal{S}'(\mathbb{R}^d)} \left[\frac{\delta}{\delta\varphi} - (m^2 - \Delta + \lambda c_\varepsilon \exp(\beta(\rho_\varepsilon * \varphi))) \right] F(\varphi) \nu(d\varphi) = 0.$$

Key idea · identify a class of ν sufficiently similar to the GFF μ^{GFF} via a coupling

$$W(\nu, \mu^{\text{GFF}}) = \inf_{\Pi} \int \|\varphi - \psi\|_*^2 \Pi(d\varphi, d\psi)$$

where $\varphi \sim \nu$, $\psi \sim \text{GFF}$ and $\varphi - \psi$ is more regular. $\Rightarrow$ Existence and uniqueness in

$$\{\nu: W(\nu, \mu^{\text{GFF}}) < \infty\}.$$

The problem is open for other models, e.g. $\Phi_{2,3}^4$ (maybe now possible via proofs of uniqueness of SQ)

what is an Euclidean field ?

There is no well defined “state space”, i.e. a good mathematical description of the local information generated by the Euclidean field.

Equivalently: we understand very poorly the family of local observables.

They do not form an algebra due to short scale singularities (this is very different from $d=1$!)

Factorisation algebras: allow to multiply only observables with disjoint support (Costello–Gwilliam). (link with OPE)

Working idea (in the subcritical regime):

Euclidean field theory is the study of deformations of the factorisation algebra of the GFF

Generalises Ito idea that Markov diffusions are locally like Brownian motion

act IV

a new class of equations

a new class of equations for Euclidean fields

Goal · identify a rigorous framework to analyse Euclidean fields

Let φ_∞ be a random field on $\mathbb{R}^d$, possibly distributional.

❶ We endow it with a decomposition over scales $(\varphi_a)_{a \geq 0}$ where φ_a is a description of φ_∞ including fluctuations at scales larger than $1/a$. $\varphi_a \rightarrow \varphi_\infty$ as $a \rightarrow \infty$ and $a \mapsto \varphi_a$ is continuous in some topology over smooth fields.

❷ Let $(\mathcal{F}_a)_a$ the filtration generated by φ_a . An **observable** is a martingale wrt. this filtration. The observable $(\mathcal{O}_a)_a$ is supported on a set $U \subseteq \mathbb{R}^d$ if

$$\mathcal{O}_a - \hat{\mathcal{O}}_a(\varphi_a, \nabla \varphi_a, \dots) \rightarrow 0, \quad \text{as } a \rightarrow \infty$$

where $\hat{\mathcal{O}}_a$ is a functional of φ_a which depends on the fields only on a $1/a$ -enlargement of the set U . A field of observables $x \in \mathbb{R}^d \mapsto (\mathcal{O}_a(x))_a$ is **local** if $\mathcal{O}_a(x)$ is supported on $\{x\}$ for all x .

E.g. if φ_∞ is a function:

$$\mathcal{O}_a(x) = \mathbb{E}[\varphi_\infty(x) | \mathcal{F}_a]$$

③ We assume that the **scale dynamics** of $(\varphi_a)_a$ is given by an Itô diffusion:

$$d\varphi_a = B_a da + dM_a, \quad d\langle M \otimes M \rangle_a = D_a^2 da$$

with adapted drift B_a and diffusivity “matrix” D_a^2 . We want that the dynamics is specified only in terms of features of φ_∞ “brought back” to the scale a . So we postulate:

- a) the existence of local observables for the “microscopic” drift $(f_a)_{a \geq 0}$ and for the “microscopic” diffusivity $(\Sigma_a^2)_{a \geq 0}$
- b) that the characteristics B_a, D_a^2 of the diffusion at scale a are given by some spatial averaging of the microscopic characteristics:

$$B_a = \dot{C}_a f_a, \quad D_a^2 = \dot{C}_a^{1/2} \Sigma_a^2 \dot{C}_a^{1/2}$$

where $(C_a)_a$ are spatial averaging operators at scale a and $\dot{C}_a = \partial_a C_a$. E.g.

$$(C_a f)(x) = \int_{\mathbb{R}^d} a^d \chi(a(x-y)) f(y) dy$$

where $\chi: \mathbb{R}^d \rightarrow \mathbb{R}$ is a mass one, positive and positive definite function with support on the unit ball. Note that we could allow also random averaging: $C_a = C_a(\varphi_a)$.

Wilson–Ito diffusions

[I. Bailleul, I. Chevyrev, MG. Wilson–Ito diffusions. arXiv (July 2023)]

Definition. A *Wilson–Ito diffusion* $(\varphi_a)_a$ is the solution of the SDE

$$d\varphi_a = \dot{C}_a f_a da + \dot{C}_a^{1/2} \Sigma_a dW_a, \quad a \geq 0$$

where W is a cylindrical Brownian motion, f_a, Σ_a^2 are local observables for the microscopic drift and diffusivity and C_a is a local averaging operator at scale a .

It describes the random field φ_∞ .

Covariant under change of scales · $A = A(a)$, $\tilde{\varphi}_a := \varphi_{A(a)}$ then $\tilde{C}_a = C_{A(a)}$ and

$$d\tilde{\varphi}_a = \dot{\tilde{C}}_a f_{A(a)} da + \dot{\tilde{C}}_a^{1/2} \Sigma_{A(a)} d\tilde{W}_a,$$

where $\tilde{W}$ is a cylindrical Brownian motion.

Trivial example · $f_a = 0$, $\Sigma_a = 1$. Then $\varphi_\infty = \int_0^\infty \dot{C}_a^{1/2} dW_a$ is a white noise, so in general the solutions are distributions.

are there non-trivial examples?

approximation strategy · fix some $A > 0$ and functional F_A and solve the forward-backward SDE:

$$d\phi_a = \dot{C}_a \mathbb{E}_a[F_A(\phi_A)] da + \dot{C}_a^{1/2} dW_a, \quad a \in (0, A).$$

where $\mathbb{E}_a = \mathbb{E}[* | \mathcal{F}_a]$. Try to send $A \rightarrow \infty$ and at the same time make F_A more and more local.

coherent germs · another approach is to “guess” the drift $f_a \approx F_a(\varphi_a)$ where F_a is the “germ”

$$f_a = F_a(\varphi_a) + R_a$$

then we have a forward–backwards system for (φ_a, R_a) :

$$\begin{cases} d\phi_a = \dot{C}_a(F_a(\varphi_a) + R_a) da + \dot{C}_a^{1/2} dW_a, \\ R_a = \mathbb{E}_a[\int_a^\infty \mathcal{L}_b F_b(\varphi_b) db + \int_a^\infty D F_b(\varphi_a) \dot{C}_b R_b db] \end{cases}$$

$$\mathcal{L}_b = \partial_b + \frac{1}{2} \Delta_{\dot{C}_b} + F_b \dot{C}_b D$$

Fully open problem in generality · I know very little about it · (some examples later)

a linear force

assume

$$f_a = \mathbb{E}_a[-A\phi_\infty] + \mathbb{E}_a[h(\phi_\infty)]$$

where A is a positive linear operator, e.g. $A = m^2 - \Delta$. Then, with $C_{\infty,a} := C_\infty - C_a$

$$\phi_\infty = \phi_a + \int_a^\infty \dot{C}_a(\mathbb{E}_a[-A\phi_\infty] + \mathbb{E}_a[h(\phi_\infty)])da + \int_a^\infty \dot{C}_a^{1/2}dW_a$$

$$\mathbb{E}_a[\phi_\infty] = \phi_a - C_{\infty,a}A\mathbb{E}_a[\phi_\infty] + C_{\infty,a}\mathbb{E}_a[h(\phi_\infty)]$$

Let $\psi_a := (1 + C_{\infty,a}A)^{-1}\phi_a$ then $\psi_\infty = \phi_\infty$ and

$$d\psi_a = \dot{Q}_a\mathbb{E}_a[h(\psi_\infty)]da + \dot{Q}_a^{1/2}dW_a, \quad \dot{Q}_a := \partial_a(A^{-1}(1 + C_{\infty,a}A)^{-1})$$

the Gaussian field

$$X_a^Q := \int_0^a \dot{Q}_c^{1/2}dW_c$$

has covariance $Q_\infty - Q_0 = (1 + A)^{-1}$, i.e. is a GFF when $A = m^2 - \Delta$.

gradient Wilson–Ito diffusions

assume now also that $h(\psi_\infty) = -DV_\infty(\psi_\infty)$, and let $(V_a)_{a \geq 0}$ be the solution to the *Polchinski equation* (HJB)

$$\partial_a V_a - \frac{1}{2} DV_a \dot{Q}_a DV_a + \frac{1}{2} \dot{Q}_a D^2 V_a = 0$$

then one can prove that

$$\mathbb{E}_a[h(\psi_\infty)] = -\mathbb{E}_a[DV_\infty(\psi_\infty)] = -DV_a(\psi_a)$$

and by performing Doob's h -transform with $d\mathbb{Q} = e^{-V_0(0) + V_a(\psi_a)} d\mathbb{P}$ also that

$$\begin{aligned} \mathbb{E}_{\mathbb{P}}[G(\psi_a)] &= \mathbb{E}_{\mathbb{Q}}[G(\psi_a) e^{V_0(0) - V_a(\psi_a)}] = \mathbb{E}_{\mathbb{P}}[G(X_a^{\mathbb{Q}}) e^{V_0(0) - V_a(X_a^{\mathbb{Q}})}] \\ &= \frac{\mathbb{E}_{\mathbb{P}}[G(X_a^{\mathbb{Q}}) e^{-V_\infty(X_\infty^{\mathbb{Q}})}]}{\mathbb{E}_{\mathbb{P}}[e^{-V_\infty(X_\infty^{\mathbb{Q}})}]} \end{aligned}$$

for any nice function G . In particular the law of ψ_∞ is given by the Gibbs measure

$$\nu_\infty(d\psi) = \frac{e^{-V_\infty(\psi)} \mu^{\mathbb{Q}_\infty}(d\psi)}{\int e^{-V_\infty(\psi)} \mu^{\mathbb{Q}_\infty}(d\psi)}.$$

Euclidean fields as Wilson–Ito fields

the class of Wilson–Itô fields comprises as a particular case the Euclidean quantum fields constructed as perturbations of a Gaussian field. They are obtained by solving **Polchinski FBSDEs** of the form

$$d\psi_a = -\dot{Q}_a \mathbb{E}_a[DV_\infty(\psi_\infty)] da + \dot{Q}_a^{1/2} dW_a.$$

optimal control formulation · Let $u_a := -\dot{Q}_a^{1/2} \mathbb{E}_a[DV_\infty(\psi_\infty)]$, test it with adapted $(v_a)_a$ and integrate:

$$\mathbb{E} \left[\int_0^\infty \langle v_a, u_a \rangle da + \left\langle \int_0^\infty \dot{Q}_a^{1/2} v_a da, DV_\infty(\psi_\infty) \right\rangle \right] = 0.$$

It is the first-order condition for the minimisation of the functional

$$\Psi(u) := \mathbb{E} \left[V_\infty(\psi_\infty^u) + \frac{1}{2} \int_0^\infty \langle u_a, u_a \rangle da \right]$$

over all adapted controls $(u_a)_{a \geq 0}$, where

$$\psi_a^u := \int_0^a \dot{Q}_b^{1/2} u_b db + \int_0^a \dot{Q}_b^{1/2} dW_b,$$

is the controlled process.

while Wilson–Ito fields are very young, we have already established some results in the same flavour by looking at FBSDE or at the stochastic control formulation of Euclidean fields.

- N. Barashkov and MG · A Variational Method for Φ_3^4 · *Duke Mathematical Journal* (2020)
- N. Barashkov and MG · The Φ_3^4 Measure via Girsanov's Theorem · *EJP* (2021)
- N. Barashkov's PhD thesis · University of Bonn (2021)
- N. Barashkov and MG · On the Variational Method for Euclidean Quantum Fields in Infinite Volume, *Prob. Math. Phys.* (2023+)
- N. Barashkov · A Stochastic Control Approach to Sine Gordon EQFT · *arXiv* (2022)
- R. Bauerschmidt, M. Hofstetter · Maximum and Coupling of the Sine-Gordon Field · *Ann. Prob.* (2022)
- F. C. De Vecchi, L. Fresta, and MG · A stochastic analysis of subcritical Euclidean fermionic field theories · *arXiv* (2022)
- N. Barashkov, T. S. Gunaratnam, M. Hofstetter · Multiscale Coupling and the Maximum of φ_2^4 Models on the Torus · *arXiv* (2023)
- R. Bauerschmidt, T. Bodineau, B. Dagallier · Stochastic Dynamics and the Polchinski Equation: An Introduction · *arXiv* (2023)
- MG and S. J. Meyer · An FBSDE for Sine–Gordon up to 6π . *arXiv* (2025)

some remarks on Wilson–Ito diffusions

- ▶ our the working hypothesis is that Wilson–Ito diffusions are natural mechanism to generate and analyse random fields
- ▶ they emerge from simple and natural assumptions and covers in principle much more than those theories that can be reached perturbatively from a Gaussian functional integral, e.g. from the path-integral picture
- ▶ they can be used for gauge theories and fields on manifolds and for Grassmann fields
- ▶ they allow for rigorous non-perturbative results in the whole space
- ▶ (hopefully) they provide a new framework for the stochastic analysis of Euclidean fields
- ▶ numerical simulations?
- ▶ still lot to understand: FBSDEs are non-trivial to analyse but PDE methods seems applicable similarly to Parisi–Wu style stochastic quantisation.

thanks

topics for other talks

global solutions & properties EQFTs

► sine-Gordon

$$V'(\varphi) = \lambda \sin(\beta \varphi)$$

► σ -models in $d=1$

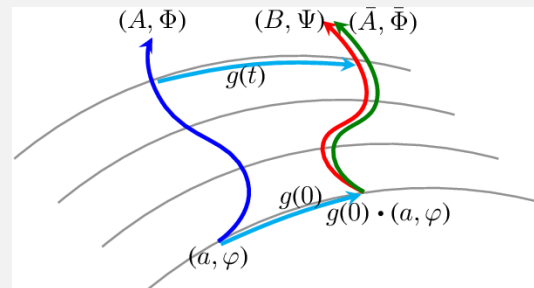
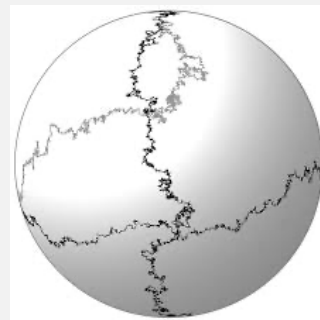
(e.g. dynamics of a loop $u: \mathbb{T} \rightarrow \mathcal{M}$ on a manifold $\mathcal{M}$)

$$\partial_t u = \Delta u + g(u)(\nabla u \otimes \nabla u) + h(u)\xi, \quad u: \mathbb{R}_+ \times \mathbb{T} \rightarrow \mathcal{M}$$

► Abelian and non-Abelian gauge theories (and Higgs)

$d=2,3$

$$\begin{cases} \partial_t A = \Delta A + g A \nabla A + g A A A + e \varphi \nabla \varphi + \xi \\ \partial_t \varphi = \Delta \varphi + e A \nabla \varphi + e A A \varphi + \lambda |\varphi|^2 \varphi + \xi \end{cases}$$



[S. Albeverio, F. C. De Vecchi, **MG** · Elliptic Stochastic Quantization · Ann. Prob. 2020 | The elliptic stochastic quantization of some two dimensional Euclidean QFTs. · *Ann. Inst. H. Poincaré, PS*, 2021]

From an idea of Parisi–Sourlas: **supersymmetric proof**

$$(m^2 - \Delta_y)\phi(y) + V'(\phi(y)) = \xi(y), \quad y \in \mathbb{R}^{d+2}$$

❶ **Change of variables** · $T(\phi) = \xi$

$$\begin{aligned} \text{Law}(\phi) &= T_*^{-1} \text{Law}(\xi) \approx \exp \left[-\frac{1}{2} \int_{\mathbb{R}^{d+2}} |(m^2 - \Delta_y)\phi(y) + V'(\phi(y))|^2 dy \right] \det(DT(\phi)) \mathcal{D}\phi \\ &\approx \exp[-\mathcal{S}] \mathcal{D}(\phi, \omega, \psi) \end{aligned}$$

$$\mathcal{S} = -\frac{1}{2} \int \omega(y)^2 dy + \int [(m^2 - \Delta_y)\phi(y) + V'(\phi(y))] \omega(y) dy + \int \bar{\psi}(z) [(m^2 - \Delta_y) + V''(\phi(y))] \psi(y) dy$$

$\psi, \bar{\psi}$ are Grassmann fields ($\psi \bar{\psi} = -\bar{\psi} \psi$) – relation with non-commutative probability

$$\mathcal{S} = -\frac{1}{2} \int \omega(y)^2 dy + \int [(m^2 - \Delta_y) \phi(y) + V'(\phi(y))] \omega(y) dy + \int \bar{\psi}(z) [(m^2 - \Delta_y) + V''(\phi(y))] \psi(y) dy$$

② The superfield

$$\Phi(Y) = \Phi(y, \theta, \bar{\theta}) = \phi(y) + \theta \bar{\psi}(y) + \bar{\theta} \psi(y) + \theta \bar{\theta} \omega(y), \quad dY = dy d\theta d\bar{\theta}, \quad Y \in \mathbb{R}^{d+2|2}$$

$$\Delta_Y = \Delta_y + \partial_\theta \partial_{\bar{\theta}},$$

$$\mathcal{S} = \frac{1}{2} \int_{\mathbb{R}^{d+2|2}} [\Phi(Y) (m^2 - \Delta_Y) \Phi(Y) + V(\Phi(Y))] dY$$

$$V(\Phi(Y)) = V(\phi(y)) + V'(\phi(y)) (\theta \bar{\psi}(y) + \bar{\theta} \psi(y) + \theta \bar{\theta} \omega(y)) - V''(\phi(y)) \bar{\psi}(y) \psi(y) \theta \bar{\theta}$$

$$\text{Law}(\phi) \approx \underbrace{\Pi_\phi}_{\text{marginal}} \exp \left[-\frac{1}{2} \int_{\mathbb{R}^{d+2|2}} [\Phi(Y) (m^2 - \Delta_Y) \Phi(Y) + V(\Phi(Y))] dY \right] \mathcal{D}\Phi$$

③ Supersymmetric localization · $z \in \mathbb{R}^2$

$$\text{Law}(\phi(z, \cdot)) \approx \exp \left[-\frac{4\pi}{2} \int_{\mathbb{R}^d} [\phi(x) (m^2 - \Delta_x) \phi(x) + V(\phi(x))] dx \right] \mathcal{D}\phi$$